

MPF1204, FISIKA KUANTUM (3 SKS)
Program Studi S2 Pendidikan Fisika



**FAKULTAS KEGURUAN DAN ILMU
PENDIDIKAN
UNIVERSITAS SEBELAS MARET (UNS)
SURAKARTA**

e Learning :

Eigen Function of Angular Momentum in The Spherical Coordinates and Hamiltonian Comuttator

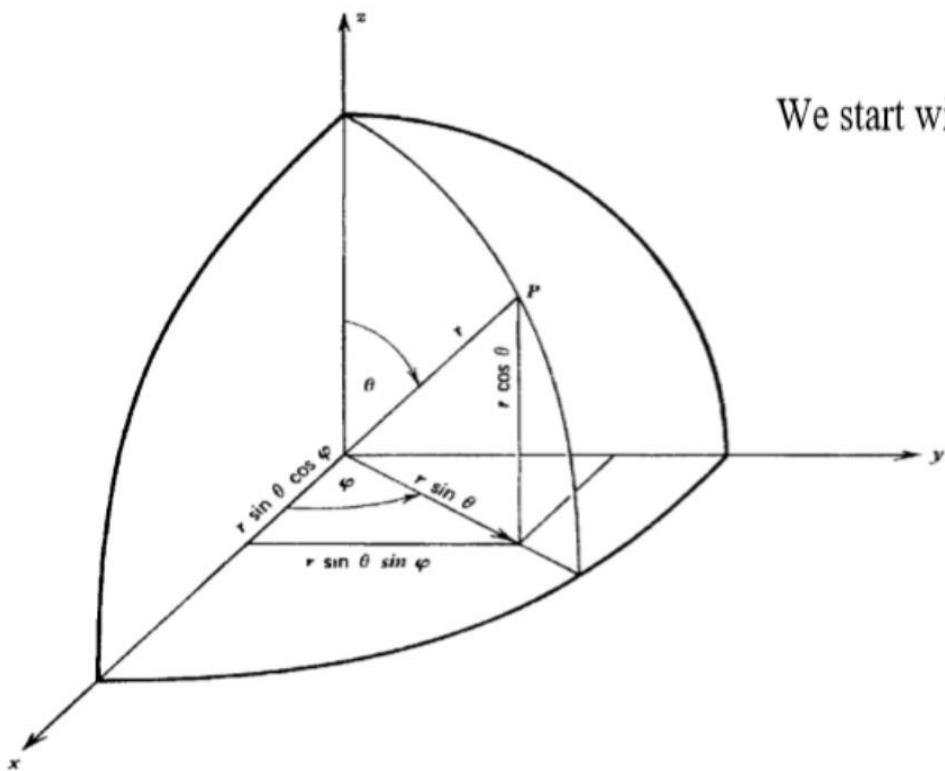
Pertemuan ke-13 :

Selasa, 12 Mei 2020 TM online Pk. 08.30 wib

DR. Suharno, M.Si



1. Eigen Function of Angular Momentum



We start with spherical coordinates, as defined in Fig. 7-2. We have

$$\begin{aligned}x &= r \sin \theta \cos \varphi \\y &= r \sin \theta \sin \varphi \\z &= r \cos \theta\end{aligned}\tag{7B-1}$$

Figure 7-2 The definition of the spherical coordinates used in the text and the relation between the cartesian coordinates (x, y, z) and the spherical coordinates (r, θ, φ) .



Eigen Function of Angular Momentum

From this it follows that

$$\begin{aligned} dx &= \sin \theta \cos \varphi dr + r \cos \theta \cos \varphi d\theta - r \sin \theta \sin \varphi d\varphi \\ dy &= \sin \theta \sin \varphi dr + r \cos \theta \sin \varphi d\theta + r \sin \theta \cos \varphi d\varphi \\ dz &= \cos \theta dr - r \sin \theta d\theta \end{aligned} \tag{7B-2}$$

These may be solved to give

$$\begin{aligned} dr &= \sin \theta \cos \varphi dx + \sin \theta \sin \varphi dy + \cos \theta dz \\ d\theta &= \frac{1}{r} (\cos \theta \cos \varphi dx + \cos \theta \sin \varphi dy - \sin \theta dz) \\ d\varphi &= \frac{1}{r \sin \theta} (-\sin \varphi dx + \cos \varphi dy) \end{aligned} \tag{7B-3}$$



Eigen Function of Angular Momentum

With the help of these we obtain

$$\begin{aligned}\frac{\partial}{\partial x} &= \frac{\partial r}{\partial x} \frac{\partial}{\partial r} + \frac{\partial \theta}{\partial x} \frac{\partial}{\partial \theta} + \frac{\partial \varphi}{\partial x} \frac{\partial}{\partial \varphi} \\ &= \sin \theta \cos \varphi \frac{\partial}{\partial r} + \frac{1}{r} \cos \theta \cos \varphi \frac{\partial}{\partial \theta} - \frac{\sin \varphi}{r \sin \theta} \frac{\partial}{\partial \varphi} \\ \frac{\partial}{\partial y} &= \sin \theta \sin \varphi \frac{\partial}{\partial r} + \frac{1}{r} \cos \theta \sin \varphi \frac{\partial}{\partial \theta} + \frac{\cos \varphi}{r \sin \theta} \frac{\partial}{\partial \varphi} \\ \frac{\partial}{\partial z} &= \cos \theta \frac{\partial}{\partial r} - \frac{1}{r} \sin \theta \frac{\partial}{\partial \theta}\end{aligned}\tag{7B-4}$$

We may use these to calculate the angular momentum operators in terms of the spherical angles. We have

$$\mathbf{L} = \mathbf{r} \times \mathbf{p}_{\text{op}} = \frac{\hbar}{i} (\mathbf{r} \times \nabla)\tag{7B-5}$$

$$L_z = \frac{\hbar}{i} \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right)$$



Eigen Function of Angular Momentum

$$\begin{aligned} &= \frac{\hbar}{i} \left[r \sin \theta \cos \varphi \left(\sin \theta \sin \varphi \frac{\partial}{\partial r} + \frac{1}{r} \cos \theta \sin \varphi \frac{\partial}{\partial \theta} + \frac{\cos \varphi}{r \sin \theta} \frac{\partial}{\partial \varphi} \right) \right. \\ &\quad \left. - r \sin \theta \sin \varphi \left(\sin \theta \cos \varphi \frac{\partial}{\partial r} + \frac{1}{r} \cos \theta \cos \varphi \frac{\partial}{\partial \theta} - \frac{\sin \varphi}{r \sin \theta} \frac{\partial}{\partial \varphi} \right) \right] \\ &= \frac{\hbar}{i} \frac{\partial}{\partial \varphi} \end{aligned} \tag{7B-6}$$

Similarly, we construct

$$\begin{aligned} L_x &= \frac{\hbar}{i} \left(y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right) \\ &= \frac{\hbar}{i} \left[r \sin \theta \sin \varphi \left(\cos \theta \frac{\partial}{\partial r} - \frac{\sin \theta}{r} \frac{\partial}{\partial \theta} \right) \right. \\ &\quad \left. - r \cos \theta \left(\sin \theta \sin \varphi \frac{\partial}{\partial r} + \frac{1}{r} \cos \theta \sin \varphi \frac{\partial}{\partial \theta} + \frac{\cos \varphi}{r \sin \theta} \frac{\partial}{\partial \varphi} \right) \right] \\ &= \frac{\hbar}{i} \left(-\sin \varphi \frac{\partial}{\partial \theta} - \cot \theta \cos \varphi \frac{\partial}{\partial \varphi} \right) \end{aligned} \tag{7B-7}$$



Eigen Function of Angular Momentum

and

$$\begin{aligned} L_y &= \frac{\hbar}{i} \left(z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right) \\ &= \frac{\hbar}{i} \left[r \cos \theta \left(\sin \theta \cos \varphi \frac{\partial}{\partial r} + \frac{1}{r} \cos \theta \cos \varphi \frac{\partial}{\partial \theta} - \frac{\sin \varphi}{r \sin \theta} \frac{\partial}{\partial \varphi} \right) \right. \\ &\quad \left. - r \sin \theta \cos \varphi \left(\cos \theta \frac{\partial}{\partial r} - \frac{\sin \theta}{r} \frac{\partial}{\partial \theta} \right) \right] \\ &= \frac{\hbar}{i} \left(\cos \varphi \frac{\partial}{\partial \theta} - \cot \theta \sin \varphi \frac{\partial}{\partial \varphi} \right) \end{aligned} \tag{7B-8}$$



Eigen Function of Angular Momentum

It is fairly straightforward to calculate

$$\begin{aligned}\mathbf{L}^2 &= L_x^2 + L_y^2 + L_z^2 \\ &\quad - \hbar^2 \left[\left(-\sin \varphi \frac{\partial}{\partial \theta} - \cot \theta \cos \varphi \frac{\partial}{\partial \varphi} \right) \left(-\sin \varphi \frac{\partial}{\partial \theta} - \cot \theta \cos \varphi \frac{\partial}{\partial \varphi} \right) \right. \\ &\quad \left. + \left(\cos \varphi \frac{\partial}{\partial \theta} - \cot \theta \sin \varphi \frac{\partial}{\partial \varphi} \right) \left(\cos \varphi \frac{\partial}{\partial \theta} - \cot \theta \sin \varphi \frac{\partial}{\partial \varphi} \right) + \frac{\partial^2}{\partial \varphi^2} \right]\end{aligned}$$

We leave it to the reader to do the algebra. The final result is

$$\mathbf{L}^2 = -\hbar^2 \left(\frac{\partial^2}{\partial \theta^2} + \cot \theta \frac{\partial}{\partial \theta} + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \varphi^2} \right) \quad (7B-9)$$

Which is just the expression in eq. (7-13).

The equation

$$\mathbf{L}^2 Y(\theta, \varphi) = \hbar^2 \lambda Y(\theta, \varphi)$$



2. Hamiltonian and Angular momentum Comutator

Consider the special case of a rotation through an angle θ about the z -axis: With

$$\begin{aligned}x' &= x \cos \theta - y \sin \theta \\y' &= x \sin \theta + y \cos \theta\end{aligned}\tag{7A-1}$$

it is easy to see that

$$r' = (x'^2 + y'^2 + z'^2)^{1/2} = (x^2 + y^2 + z^2)^{1/2} = r\tag{7A-2}$$

and

$$\begin{aligned}\left(\frac{\partial}{\partial x'}\right)^2 + \left(\frac{\partial}{\partial y'}\right)^2 &= \left(\cos \theta \frac{\partial}{\partial x} - \sin \theta \frac{\partial}{\partial y}\right)^2 + \left(\sin \theta \frac{\partial}{\partial x} + \cos \theta \frac{\partial}{\partial y}\right)^2 \\&= \left(\frac{\partial}{\partial x}\right)^2 + \left(\frac{\partial}{\partial y}\right)^2\end{aligned}\tag{7A-3}$$



Hamiltonian and Angular momentum Comutator

Since the Hamiltonian has an invariance property, we expect a conservation law, as we saw in the case of parity. To identify the operators that commute with H , let us consider an infinitesimal rotation about the z -axis. Keeping terms of order θ only so that

$$\begin{aligned}x' &= x - \theta y \\y' &= y + \theta x\end{aligned}\tag{7A-4}$$

we require that

$$Hu_E(x - \theta y, y + \theta x, z) = Eu_E(x - \theta y, y + \theta x, z)\tag{7A-5}$$

If we expand this to first order in θ and subtract from it

$$Hu_E(x, y, z) = Eu_E(x, y, z)\tag{7A-6}$$

we obtain from the term linear in θ

$$H\left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x}\right)u_E(x, y, z) = E\left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x}\right)u_E(x, y, z)\tag{7A-7}$$



Hamiltonian and Angular momentum Comutator

The right side of this equation may be written as

$$\left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right) E u_E(x, y, z) = \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right) H u_E(x, y, z) \quad (7A-8)$$

If we define

$$L_z = \frac{\hbar}{i} \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right) = x p_y - y p_x \quad (7A-9)$$

then (7A-7) and (7A-8) together read

$$(H L_z - L_z H) u_E(x, y, z) = 0$$

Since the $u_E(\mathbf{r})$ form a complete set, this implies the operator relation

$$[H, L_z] = 0 \quad (7A-10)$$



Hamiltonian and Angular momentum Comutator

holds. L_z is the z -component of the operator

$$\mathbf{L} = \mathbf{r} \times \mathbf{p} \quad (7A-11)$$

which is the angular momentum. Had we taken rotations about the x - and y -axis, we would have found, in addition, that

$$\begin{aligned} [H, L_x] &= 0 \\ [H, L_y] &= 0 \end{aligned} \quad (7A-12)$$

Thus the three components of the angular momentum operators commute with the Hamiltonian; that is, the angular momentum is a constant of the motion. This parallels the classical result that central forces imply conservation of the angular momentum.



Exercises :

1. Calculate $\langle l, m_1 | L_x | l, m_2 \rangle$ and $\langle l, m_1 | L_y | l, m_2 \rangle$.
2. Calculate $\langle l, m_1 | L_x^2 | l, m_2 \rangle$ and $\langle l, m_1 | L_y^2 | l, m_2 \rangle$. (*Hint:* Things will be easier with the use of L_{\pm} , L_{\mp}^2 , and so on.)



Thank you